

Mathematical Model of Biological Tissues towards Morphogenesis

Mr. Saroj Sarkar
Research Scholar
sarojsarkar89@gmail.com

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Abstract

Lander, Nie, and Wan provided a basic mathematical model of morphogenesis distribution (signalling molecules responsible for cell differentiation and tissue pattern formation) in 2002. The model is made up of two equations: a parabolic PDE with a dynamic boundary condition that models the distribution of free morphogenesis, and an ODE that describes the evolution of bound receptors. Diffusion, degradation, and reversible binding are three biological processes that are considered. We show that solutions exist, are unique, and have asymptotic behaviour.

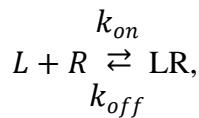
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1. Introduction

Babylonians and Egyptians are known to have lived before 300 BC, but before BC, what the whole line looks like in a larger segment, the larger ones are smaller. When this concept revealed a simple ratio used by the sculptor Fidia more than two centuries ago. Parthenon. The idea of using such simple geometric rules to decipher the harmony and symmetry behind the various patterns that are regularly observed in nature still fascinates people today and against the numbers of heaven. Justifies medieval beliefs. The best, and the form is pre-ordered as one organization of matter resulting from the sacred act of creation. In contrast to such preformationist views, Aristotle proposed the so-called epigenetic hypothesis in his books *Hystoria Animalium* and *De Generatione Animalium* (c. 350 BC). There was no life [6]. Although systematically far from the rigorous current scientific methodologies, such publications are primarily based on teleological principles and have recorded a surprising amount of biological observations of embryos.

2. The mathematical model.

A few researchers have proposed a few morphogenetic scattering speculations throughout the past 10 years. In the following part, we'll take a gander at a basic numerical model introduced via Lander, Nie, and Wan [6]. They consider the advancement of decapentapresia (Dpp) as one of the morphogenetic processes in the Drosophila larval wing circle. Utilizing the one-layered locale (0), the model works on the state of the wing circle. Morphogen Dpp (linand) is shown by L, receptors per unit of extracellular space are demonstrated by R, complex ligand receptors are shown by LR, and their separate focuses are [L], [R], and [LR]. The accompanying recipe sums up the cycle that happens when ligand L ties to receptor R to frame complex LR as well as the other way around.



The binding and dissociation rate constants are k_{on} and k_{off} , respectively. We assume that the number of free and bound receptors remains constant over time.

$$[R] = R_{tot} - [LR],$$

R_{tot} addresses the absolute receptor focus per unit of extracellular space. Presumption [1] works on the issue by lessening the quantity of conditions. Debasement of [LR] ought to be considered throughout an extensive stretch of time (see Lander, Nie, Vargas, Wan [8]). As made sense of in the past segment, we accept a straight dispersion of [L] with dissemination consistent of d.

Where

$$\frac{\partial}{\partial t} [L] - d \frac{\partial^2}{\partial x^2} [L] = - k_{on} R_{tot} [L] + k_{on} [L][LR] + k_{off} [LR], \quad x > 0, t > 0, \quad (2)$$

$$d \sim 10^{-13} M^2 s^{-1}; k_{on}, R_{tot} \sim 1 - 10^{-2} s^{-1};$$

$$k_{on} \sim 10^5 M^{-1} s^{-1}; k_{off} \sim 10^{-6} - 10^{-8} s^{-1}$$

3. An historical overview of morphogenetic theories

The motivation behind this segment is to give a timetable outline of the main occasions in development and natural revelations, from old to present and current methodologies..

3.1 Epigenesist versus pre-formations: from ancient times to the advent of microscopy.

Babylonians and Egyptians are known to have lived before 300 BC, but before BC, what the whole line looks like in a larger segment, the larger ones are smaller. When this concept revealed a simple ratio used by the sculptor Fidia more than two centuries ago. Parthenon. The idea of using such simple geometric rules to decipher the harmony and symmetry behind the various patterns that are regularly observed in nature still fascinates people today and against the numbers of heaven. Justifies medieval beliefs. The best, and the form is pre-ordered as one association of issue coming about because of the consecrated demonstration of creation. Rather than such preformationist sees, Aristotle proposed the alleged epigenetic theory in his books Hystoria Animalium and DeGeneratione Animalium (c. 350 BC). There was no life [6]. Albeit strategically incredibly distant from a thorough current logical philosophy, such works were for the most part founded on teleological originations, yet they archived a shocking measure of organic discoveries on undeveloped organisms.

3.2 The birth of modern embryology: evolutionary theories and mechanical causation

Darwinism and the birth of genetics dominated the nineteenth century, but so did the belief that plants and creatures are made up of cells and that cell division governs reproduction, according to Rudolf Virchow's concept omnis cellulae ex cellula [116]. Furthermore, advancements in microscopy were crucial in the development of modern experimental embryology. Ernst Haeckel proposed the summarization hypothesis, which asserts that the formative phases of undeveloped development are an impression of their ancestors' grown-up advancement [55]. The ontogeny (i.e., the advancement of a singular creature) restates phylogeny, as per this biogenetic regulation (for example the advancement of an animal groups). Regardless of being immediately exposed by Wilhelm His, who exhibited with tests that stomach tube morphogenesis could be addressed utilizing a mechanical causation standard [93], this biogenetic speculation continued in the logical world until Wilhelm Roux's examinations close to the furthest limit of the nineteenth 100 years. He moved the focal point of formative science from advancement (ie, a definitive objective) to the fundamental system (ie, cause) and conducted several studies on embryos to investigate self-differentiation. [100] Lou influences a new method of experimental developmental biology, with pioneering data on the regulation or ability to develop normally even when part of the embryo is removed, and the induction or effect on another development of one cell or tissue. Provided. 56]. Interestingly, it was in 1924 that Spemann and Mangold exhibited the latter phenomenon when part of the embryonic tissue of one amphibian was transplanted into another amphibian embryo. [107] They found that the transplanted tissue

caused partial second embryonic development: the host tissue induced the growth of the host tissue.

3.3 The chemical bases of morphogenesis

Rapid advances in biochemical experimental techniques have led to increasing attention to the chemical and molecular mechanisms that regulate Mor.

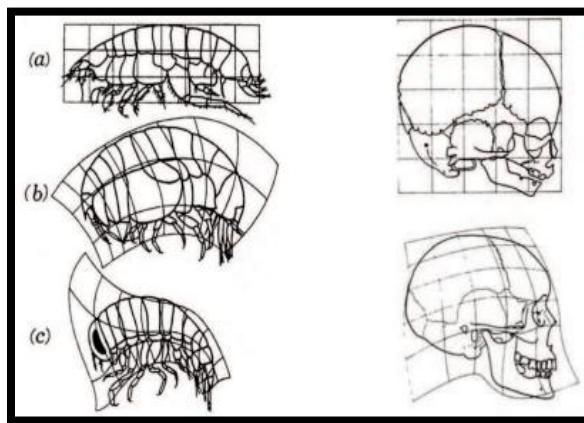


Fig. 1 Left: Following [29], the transformation grid was used to reshape the small amphipods (a) *Harpinia plumosa* to the shape of two other species of the same family (b) *Stegocephalus inflatus* and (c) *Hyperia Galba*.

Phylogenetic Processes The formation of Phylogenetic Processes Patterns, or the emergence of ordered and living structures in space and time, is the subject of a new field of applied mathematics. Alan Turing's study on the chemical basis of morphogenesis [114], published in 1952, is an original study in biological mathematics. It was his only contribution to the field of morphogenesis, but it was an ingenious effort to show how mathematical theory could predict biological discoveries. Turing has developed a reaction-diffusion model in which at least two chemical species in a biomaterial undergo a chemical reaction. He named them morphogens to emphasize their function in the formation of new patterns. They are biological compounds that give embryonic tissue specific shape control capabilities and are similar to the evocators proposed by Waddington [117]. The biological system is in a steady state, as defined by the uniform concentration of the two reactants in the absence of diffusion. Instability occurs when the reaction and diffusion parameters are adjusted to a range of values and a stable, non-uniform pattern appears. This was an unexpected conclusion, as diffusion was thought to cause the system to become chaotic rather than provide an ordered pattern.

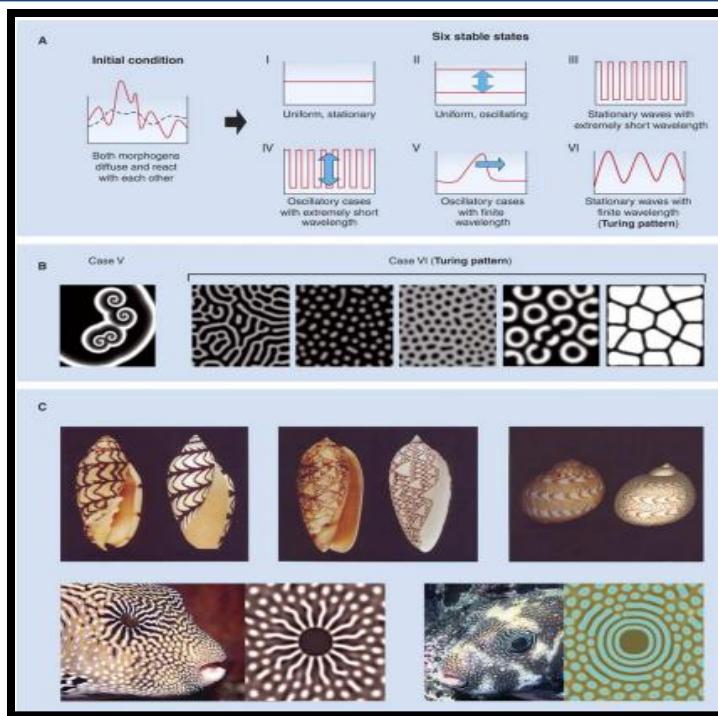


Fig. 2 Turing's response dispersion model: (a) Examples of six stable-state arrangements of the Turing model. (B) In the supposed Turing design, standing waves with a restricted frequency are shaped.

This is displayed as Case VI. (C) Reproduction of the organic example shaped by the altered response dissemination framework. [73] was utilized as a beginning stage.

The so-called short-range activation, long-range suppression [48] mechanism is triggered when the diffusion coefficients of the two morphogens differ significantly. It is believed that the two morphogens act as activators and inhibitors, respectively, and can interact with each other. A slight change in uniform concentration can cause an increase in activator concentration, cause feedback, and produce one of the Turing patterns (see). **Figure 2**.

4. A continuous chemo-mechanical theory of morphogenesis

The thermodynamics of open frameworks give a characteristic structure to the perceptible numerical depiction of the morphogenesis of living life forms. From now on, it is called continuous distribution of substances. As a result, we ignore the microscopic dynamics of the underlying biological phenomenon that functions intracellularly or at the cellular level, and instead use a field averaged at the tissue level to represent the required chemical-mechanical properties. The principle of physical balance that governs the morphogenic process of an organism as a function of both mechanical and chemical inputs is derived in the next section using some basic kinematic terms.

4.1 Basic kinematic notions

Let B_0 and B_a be two regions occupied by organisms at two different time points. Here, ER3 represents a three-dimensional Euclidean space. A map that transforms the tissue from the initial B_0 configuration to the final B_a configuration and can be characterized as a C^1 diffeomorphism:

$$\chi : \mathcal{B}_0 \rightarrow \mathcal{B}_a \quad \mathbf{x} = \chi(\mathbf{X}), \quad \text{with inverse} \quad \mathbf{X} = \chi^{-1}(\mathbf{x})$$

4.2 Balance of mass

First, since the morphogenesis process usually involves growth and remodeling, let us calculate the mass balance of biomaterials as a function of biochemical variables [40,26]. In spatial coordinates, the global form of the mass balance of the evolving object X is:

$$\frac{d}{dt} \int_{\mathcal{B}_a} \rho \, dv = \int_{\mathcal{B}_a} \omega \, dv$$

Where ρ is the total material time derivative and d / dt is the spatial mass density. The result is the current internal mass production rate per unit of volume, and for the sake of brevity, the non-convective mass flow rate is ignored. Interested readers should refer to [39] for an introduction to a large number of self-propagating words. Mass production is a measure of cell proliferation within a material and can be affected by biochemical factors such as growth signals, nutrition and morphogens. If we use $c_i(x, t)$ to represent the concentration of a common i -th species per unit volume, we can assume that it is equal to (c_i, F) . In addition, because we are dealing with signals or low mass molecules, such species can be treated as low inertia internal variables [21] and the following equilibrium rules can be applied.

$$\frac{d}{dt} c_i - \text{div } \mathbf{J}_{ci} = -\xi_i(\mathbf{F}, \text{Grad } \mathbf{F})$$

4.3 Balance of internal energy and entropy inequality

Utilizing the [40] approach, the inner energy per unit mass is addressed by a scalar capacity that mirrors the wellspring of the inside energy related with the I -th solute c_i . Subsequently, the main law of thermodynamics can be communicated in the worldwide adaptation as:

$$\frac{d}{dt} \int_{\mathcal{B}_a} \rho \varepsilon \, dv = \int_{\mathcal{B}_a} (\omega \varepsilon - \xi_i c_i \mu_i + r_0 + \text{tr}(\sigma \mathbf{d})) \, dv + \int_{\partial \mathcal{B}_a} \mathbf{n} \cdot (\mu_i \mathbf{J}_{ci} - \mathbf{Q}) \, ds$$

$d = (\text{grad}v + (\text{grad}v) T)/2$, where r_0 is the outer hotness input per unit volume, Q is the hotness stream, I is the substance capability of the I -th kind, and I is the expansion in inside energy. A continuum because of the actual assimilation of diffuse synthetics. Irreversible terms and temperature inclination conditions are disregarded for straightforwardness. By embedding an articulation. Condition (24) (10,12,15):

$$\rho \dot{\epsilon} = tr(\sigma d) + \mu_i \dot{c}_i - \text{div}Q + r_0 + J_{ci} \cdot \text{grad} \mu_i$$

4.4 Balance laws for interfacial morphogenetic processes

The morphogenetic cycle can happen on different scales, every one of not entirely settled by the typical length that portrays the neighborhood communication between the large scale manufacturing of biomolecules and the response dispersion properties. For instance, cells replicate within the confined zones created by the morphogenetic signaling diffusion front, and macroscopic remodeling occurs at the tissue level to achieve homeostasis [77]. As a result, the accounting law obtained in the previous section can be applied when the mass change is localized to a small amount. Consider a biological system composed of two different materials that occupy extended adjacent domains $V(t)$ and $V+(t)$ separated by a moving surface (t) with an external normal n . Such surfaces are treated as immaterial interfaces with thermomechanical properties [28] and their spatial positions are $x = x(u_1, u_2)$, tangent-based $a_l = x, u_l$, and $l = (1,2)$. The surface parametric velocity v can be decomposed as follows.

$$\bar{v}_\Sigma = \bar{v}_{\Sigma s} + \bar{v}_{\Sigma n} n_\Sigma$$

5. Free-boundary morphogenesis for fluid-like living matter

In the past area, we continued with the overall hypothesis of compound mechanics of morphogenesis. This requires the particular of constitutive conditions that portray the natural rheological properties of life forms, from thick fluids to delicate solids. We should begin with [18]. This is a straightforward free limit morphogenesis issue in a fluid like organic framework. Zeroing in on chemotactic extension of a gathering of living cells demonstrated as an incompressible Newtonian liquid, the it are gotten to follow results.

$$\sigma = -pI + \eta d$$

Where p is the viscosity of the fluid and v is the hydrostatic pressure. The NavierStokes equation governs motion as in equation (1). (15):

$$\rho \frac{d\mathbf{v}}{dt} = \rho \mathbf{b}_v - \text{grad}p + \eta \nabla^2 \mathbf{v}$$

Where ∇^2 indicates the spatial Laplacian and indicates the constant spatial density of the cell. This is about the same as the density of water.

5.1 Definition of the chemotactic model in a Hele Shaw cell

We consider a gathering of live cells contained between two equal level plates isolated by a short hole of length l and involving the space $x \times b$ of a Hele Shaw cell, where b is the rectilinear boundary. As found in Figure 5, the cells are lined by an inviscid liquid that possesses the locale $x > b$ (a). Let $z = 0$ signify the plane between the plates, and L mean the trademark plainly visible size of the morphogenetic design we're searching for. Thus, the speed part in z is considerably not exactly the other two, which not entirely settled by the accompanying scale [54]:

$$\frac{\partial^2 v_i}{\partial x^2} \ll \frac{\partial^2 v_i}{\partial z^2}$$

The velocity component v_i is represented by the index $i = (x, y)$. The normal cell doubling period is about 1 day, and because the fluid is so viscous, the associated flow occurs at a low Reynolds number, and the $2v$ term associated with convection and gravity becomes dominant, equation (1). It can be obtained. As (43):

$$\text{grad } p = \eta \frac{\partial^2 v_i}{\partial z^2}$$

5.2 Dimensionless form of the governing equations

The following values can be used to recast the governing equation in a dimensionless form:

$$t_c = \gamma_n^{-1}; \quad L_c = \sqrt{\frac{D_c}{\gamma_c}}; \quad v_c = \sqrt{D_c \gamma_c}; \\ p_c = \frac{p_c}{k_p}; \quad c_c = c(x \rightarrow \infty)$$

Where t_c , L_c , v_c , p_c , c_c show the trademark time, length, speed, tension, and synthetic focus. By characterizing dimensionless factors for such trademark values, the equations (48, 15, 51) can be rewritten in dimensionless form:

$$\dot{\bar{c}} = \begin{cases} \bar{\nabla}^2 \bar{c} - \bar{c} & \text{if } \bar{x} \leq \bar{x}_b \\ \bar{\nabla}^2 \bar{c} & \text{if } \bar{x} > \bar{x}_b \end{cases} \\ \bar{v} = -\bar{\text{grad}} \bar{p} \\ \bar{\nabla}^2 \bar{p} = -\bar{\nabla}^2 \bar{n}; \quad \text{with} \quad G = \frac{K_c c_c}{\rho D_n}$$

Bars indicate variables that are not dimensioned. G in Eq. (58) shows the relationship between morphogen-induced mass production and the mass emission rate caused by the viscous process.,

is the only dimensionless parameter that results. Furthermore, Eqs. (52-54)'s boundary conditions rewrite as:

$$\bar{p} = \bar{p}_0 - \delta \bar{C}; \text{ with } \delta = \frac{\sigma K_p}{D_n} \sqrt{\frac{\gamma_n}{D_n}} = \frac{\sigma \ell^2}{8\eta v_c L_c^2} \quad \text{at } \bar{x} = \bar{x}_b$$

$$\frac{d\bar{x}_b}{d\bar{t}} \cdot n_b = \bar{v}(x_b) \cdot n_b$$

$$\bar{c}(\bar{x}_b) = \bar{c}(\bar{x}_b^+) ; \bar{\nabla} \bar{c}(\bar{x}_b^-) \cdot c_b = \bar{\nabla} \bar{c}(\bar{x}_b^+) \cdot n_b$$

The mathematical model generated by equation (5661) is fully driven by the dimensionless parameter G and specified by equations (58, 59) that govern the evolution of the biological system during morphogenesis., according to the Buckingham theorem [13]. The ratio of surface tension to viscous forces is represented by in physical terms. In the next analysis, we'll remove the bars to make the notation more succinct.

5.3 Pattern formation in the nonlinear regime

The change plot of the direct soundness examination showed that the limits of the straight line are generally shaky at long frequencies. The scattering bend is like that saw in other hydrodynamic issues, for example, the shakiness of Saffmann-Taylor [101] and Mullins-Sekerka [89]. Comparatively, moving fronts are expected to produce rapidly expanding waves whose shape is determined by a completely non-linear effect. Numerical studies of the formation of fingering instability when expanding the free straight plane of the cell monolayer, and some of the epithelial cells showing the formation of fingering instability when expanding the free straight plane of the cell monolayer Explain the experimental results of [94]. He transformed the problem into a finite element code and proposed a computational solution [49]. The equations (56, 58) were solved on a triangular grid and the moving interface was adjusted at each iteration. In summary, for a particular nutrient concentration at time t_m , first set the pressure p_m according to equation (1). (58) Next, the velocity field is calculated using Darcy's law. Therefore, you can explicitly change the limits using the implicit Euler approach. (56) The concentration at time $t_m + 1$ has been resolved. To ensure the stability of the semi-implicit algorithm at each iteration

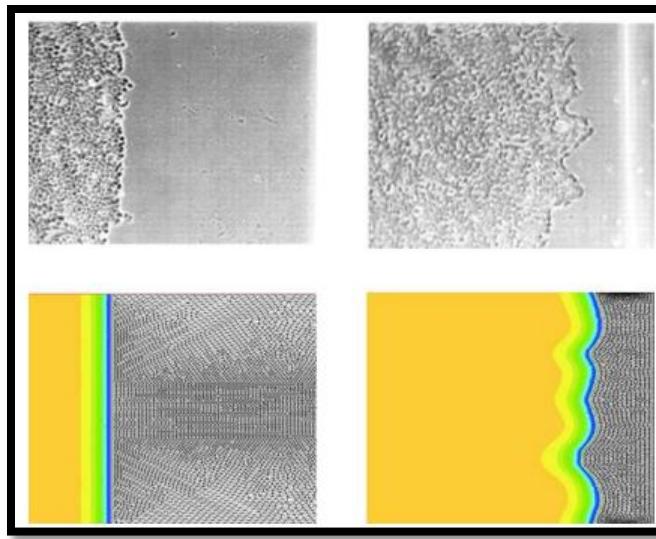


Fig.3 Monolayer growth of tumor cells over 24 hours in a wound healing experiment. [108] was used as a starting point. (B) Numerical simulation of the proposed infinite morphogenetic model in complete nonlinearity.

At the m th iteration, the time step $t = (tm + 1tm)$ is equal to the ratio of the minimum mesh size to the maximum velocity factor factor. Finally, a continuous piecewise quadratic P2 Lagrangian component is used to perform spatial discretization with finite components. Since the cross section size influences the littlest recognizable unsound frequency, we adaptively changed the lattice to change how much components on the shape of the state at each time step..

6. Conclusion

In this segment, we have encouraged a constant manufactured mechanical method for managing morphogenesis by cultivating the harmony principles and formative laws of volume and mark of communication processes. Using both theoretical procedures and reenactment gadgets, the proposed theory was applied to the examination of model age for both fluid like natural structure models and solid like normal system models. .. Organic materials, then again, have different rheological properties during such confined ideal way of behaving [115]. Furthermore, the morphogenesis interaction might incorporate microstructure revamp cycles, for example, B. Cell replication and/or relocation causes liquid like pressure unwinding peculiarities in the day to day course of events [96, 97]. To catch the infinitesimal cycle down to molecule cell size, a constant methodology should be joined with the technique for individual cell-based models [35]. At long last, really demonstrating work should be done to concentrate on the job of morphogens in development, shape, and size guideline. From one perspective, their neighborhood fixation can prompt an expansion in volume because of arbitrary cell multiplication. Need direction can happen during cell division and rivalry between various cell populaces. Then again, more spatial

arrangement is expected to change arbitrary development into uniform development. Uniform development ought to ultimately stop when it arrives at the right size. Accordingly, shape control can be impacted by both cell overflow and by large size, implying that the spatial slope of morphogen can give cells a layered acknowledgment instrument. The disclosure [125] that phones can evaluate inclinations by contrasting their flagging levels with those of nearby cells through particular administrative pathways upheld this early idea.

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